

Gauge-Origin Independent Density-Functional Theory Calculations of Vibrational Raman Optical Activity

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We present density-functional theory calculations of vibrational Raman optical activity. Gauge-origin independence of the results is ensured by using London atomic orbitals, and the frequency dependence is included by using analytical response theory. The Raman optical activity circular intensity differences are obtained by combining analytical calculations of linear response functions with numerical differentiation with respect to nuclear distortions. The results are compared with experiment, previous Hartree–Fock calculations, and a more recent approximate sum-over-states approach.

1. Introduction

The applicability of the Raman technique in solution combined with the stereochemical information provided by the sign of the circular vibrational Raman scattering intensity difference gives vibrational Raman optical activity (VROA) a great potential for structural studies of biomolecules.^{1–3} Nevertheless, since its discovery in the 1970s, VROA has not attracted much interest, mainly due to the lack of a commercially available spectrometer. Indeed, the first announcement of a commercially available VROA spectrometer was made only last year.⁴

Although the vibrational frequencies observed in a VROA spectrum can be assigned in much the same way as for an ordinary Raman spectrum, the sign of the circular intensity difference (CID) is difficult to predict a priori because of the lack of empirical rules for correlating the sign of the intensity to the absolute stereochemistry of the molecule. It can be expected that empirical rules for the CID sign in relation to stereochemistry will receive more attention as VROA becomes more widespread; currently, the use of theoretical methods for determining the sign of the CIDs in relation to the absolute stereochemistry of the molecule is needed for an unambiguous assignment.⁵

For several reasons, the theoretical calculation of VROA CIDs is a challenging task. First, because the property vanishes for a static electromagnetic field, either linear response methods^{6,7} or the static-limit approximation introduced by Amos⁸ must be used. Second, the tensors contributing to the VROA intensities are third-order properties, requiring so far the use of an expensive numerical differentiation.^{7,9} Third, one of the contributing tensors involves the interaction with the magnetic component of the electromagnetic field, making approximate theoretical calculations gauge-origin dependent, this problem can be solved using London atomic orbitals.^{7,10–12} Fourth, it is well-known that Hartree–Fock harmonic force fields are of

rather poor quality,^{13,14} which might affect the intensities because the tensors are transformed to the normal coordinate basis.^{7,15}

Only a few implementations of the full VROA CIDs have been presented,^{7,9} and the literature on calculations of VROA CIDs is still rather limited—see ref 5 for a review. Recently, Bouř and co-workers introduced an approximate sum-over-states (SOS) scheme that significantly reduces the time required for calculating VROA intensities.^{16,17} These schemes are sometimes combined with an approximate evaluation of some of the generalized polarizabilities involved in the VROA CID, known as the polar model,¹⁸ which we will denote as SOS+Polar. Their scheme has also used more accurate density-functional theory (DFT) force fields.¹⁷ In the SOS approach, the summation over states is truncated, thus introducing errors in the intensities, although the CID signs probably remain correct. However, in evaluating the performance of this approximate method, the CIDs obtained at the DFT level were compared with intensities calculated with a mixed numerical-analytical scheme at the Hartree–Fock level. This is unfortunate, as it becomes difficult to disentangle the errors arising from the SOS truncation from the errors inherent in the Hartree–Fock approximation.

In this paper, we present a mixed numerical-analytical scheme for calculating VROA intensities consistently by DFT. Gauge-origin independence is ensured by the use of the London atomic orbitals¹⁰ and the dependence on the frequency is accounted for by linear response theory.^{6,7} The derivatives of the different generalized polarizabilities with respect to nuclear distortions are obtained by numerical differentiation of the analytically calculated polarizabilities. This work therefore combines our previous Hartree–Fock implementation of VROA⁷ with our recent implementation of gauge-origin independent optical rotation at the DFT level.¹⁹

The paper is organized as follows. In section 2, we give a brief summary of the theory for VROA calculations. In section 3, we provide some technical details about the calculations presented in this paper. Section 4 summarizes our results and compares these to experimental observations and earlier theoretical works. Finally, we give some concluding remarks in section 5.

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2. Theory

The theoretical foundation of VROA was developed by Barron and Buckingham,²⁰ following the theoretical prediction of the scattering mechanism by Atkins and Barron.²¹ In particular, it was demonstrated that the VROA intensity is determined by the geometry dependence of three generalized polarizabilities—namely, the ordinary electric dipole–electric dipole polarizability

$$\alpha_{\alpha\beta} = 2 \sum_{n \neq 0} \omega_{n0} \frac{\text{Re}[\langle 0 | \mu_{\alpha} | n \rangle \langle n | \mu_{\beta} | 0 \rangle]}{\omega_{n0}^2 - \omega^2} \quad (1)$$

the mixed electric dipole–magnetic dipole polarizability

$$G'_{\alpha\beta} = -2 \sum_{n \neq 0} \omega \frac{\text{Im}[\langle 0 | \mu_{\alpha} | n \rangle \langle n | m_{\beta} | 0 \rangle]}{\omega_{n0}^2 - \omega^2} \quad (2)$$

and the mixed electric dipole–electric quadrupole polarizability

$$A_{\alpha\beta\gamma} = 2 \sum_{n \neq 0} \omega_{n0} \frac{\text{Re}[\langle 0 | \mu_{\alpha} | n \rangle \langle n | \Theta_{\beta\gamma} | 0 \rangle]}{\omega_{n0}^2 - \omega^2} \quad (3)$$

Here ω_{n0} is the excitation energy from the ground state $|0\rangle$ to the excited state $|n\rangle$ and ω is the frequency of the incident light. The summations are over all the (exact) electronically excited states of the molecule. In these equations, we have also introduced the electric dipole operator (atomic units used throughout)

$$\mu_{\alpha} = - \sum_i r_{i\alpha} \quad (4)$$

the magnetic dipole operator

$$m_{\alpha} = - \frac{1}{2} \sum_i l_{i\alpha} \quad (5)$$

and the electric quadrupole operator

$$\Theta_{\alpha\beta} = - \frac{1}{2} \sum_i (3r_{i\alpha}r_{i\beta} - r_i^2 \delta_{\alpha\beta}) \quad (6)$$

The summations are over all electrons, and r_i and l_i are the position and angular momentum of electron i , respectively.

Of particular interest for the calculation of VROA CIDs is the mixed electric dipole–magnetic dipole polarizability in eq 2. The magnetic dipole moment depends on the choice of gauge origin, leading to an unphysical dependence on this origin in approximate calculations. However, we have previously demonstrated that, for variational wave functions such as Hartree–Fock and multiconfigurational self-consistent field (MCSCF) wave functions,⁷ the exact origin dependence is recovered even for finite basis sets, provided London atomic orbitals are used.^{10,22} Note that, unlike its trace, the individual components of the electric dipole–magnetic dipole polarizability tensor are not gauge-origin independent, even for the exact wave function.

The generalized polarizability eq 2 appears also in calculations of optical rotation.²³ Stephens and co-workers have demonstrated how the gauge-origin independent DFT calculation of optical rotation is achieved by the use of London atomic orbitals^{7,10} combined with time-dependent DFT. We recently implemented¹⁹

this formalism in the Dalton program system.²⁴ To use this implementation of London-orbital DFT optical rotation¹⁹ for the calculation of VROA CIDs, we note that the generalized polarizabilities in eqs 1–3 can be formulated as linear response functions⁶

$$\alpha_{\alpha\beta} = - \langle \langle \mu_{\alpha}; \mu_{\beta} \rangle \rangle \quad (7)$$

$$G'_{\alpha\beta} = -i \langle \langle \mu_{\alpha}; m_{\beta} \rangle \rangle_{\omega} \quad (8)$$

$$A_{\alpha\beta\gamma} = - \langle \langle \mu_{\alpha}; \Theta_{\beta\gamma} \rangle \rangle_{\omega} \quad (9)$$

Of particular interest here is the CID introduced by Barron and Buckingham²⁰

$$\Delta_{\alpha} = \frac{I_{\alpha}^R - I_{\alpha}^L}{I_{\alpha}^R + I_{\alpha}^L} \quad (10)$$

where I_{α}^R and I_{α}^L are the scattered intensities with α polarization for right and left circularly polarized light, respectively. More specifically, the CIDs for right-angle and backscattering are given as

$$\Delta_z(90^\circ) = \frac{6\beta(G')^2 - 2\beta(A)^2}{3c\beta(\alpha)^2} \quad (11)$$

$$\Delta_z(180^\circ) = \frac{24\beta(G')^2 + 8\beta(A)^2}{3c\beta(\alpha)^2} \quad (12)$$

where we have expressed the CIDs in terms of the anisotropic invariants

$$\beta(\alpha)^2 = \frac{1}{2}(3\alpha_{\alpha\beta}\alpha_{\alpha\beta} - \alpha_{\alpha\alpha}\alpha_{\beta\beta}) \quad (13)$$

$$\beta(G')^2 = \frac{1}{2}(3\alpha_{\alpha\beta}G'_{\alpha\beta} - \alpha_{\alpha\alpha}G'_{\beta\beta}) \quad (14)$$

$$\beta(A)^2 = \frac{1}{2}\omega\alpha_{\alpha\beta}\epsilon_{\alpha\gamma\delta}A_{\gamma\delta\beta} \quad (15)$$

using implicit summation over repeated indices, $\epsilon_{\alpha\beta\gamma}$ is the unit third-rank antisymmetric tensor, and c is the speed of light.

For the vibrational Raman scattering process, we do not need the generalized polarizabilities eqs 1–3 but rather the corresponding vibrational transition moments. We do not repeat the derivations of ref 25 but note that, by carrying out a perturbation expansion of the generalized polarizability tensors and the vibrational wave function, and by invoking the Placzek approximation,²⁶ we can approximate the vibrational transition moments as

$$\langle \nu_0 | \alpha_{\alpha\beta} | \nu_{1p} \rangle \langle \nu_{1p} | \alpha_{\alpha\beta} | \nu_0 \rangle = \frac{1}{2\omega_p} \left(\frac{\partial \alpha_{\alpha\beta}}{\partial Q_p} \right) \Big|_{r_e} \left(\frac{\partial \alpha_{\alpha\beta}}{\partial Q_p} \right) \Big|_{r_e} \quad (16)$$

$$\langle \nu_0 | \alpha_{\alpha\beta} | \nu_{1p} \rangle \langle \nu_{1p} | G'_{\alpha\beta} | \nu_0 \rangle = \frac{1}{2\omega_p} \left(\frac{\partial \alpha_{\alpha\beta}}{\partial Q_p} \right) \Big|_{r_e} \left(\frac{\partial G'_{\alpha\beta}}{\partial Q_p} \right) \Big|_{r_e} \quad (17)$$

$$\langle \nu_0 | \alpha_{\alpha\beta} | \nu_{1p} \rangle \langle \nu_{1p} | \epsilon_{\alpha\gamma\delta} A_{\gamma\delta\beta} | \nu_0 \rangle = \frac{1}{2\omega_p} \left(\frac{\partial \alpha_{\alpha\beta}}{\partial Q_p} \right) \Big|_{r_e} \epsilon_{\alpha\gamma\delta} \left(\frac{\partial A_{\gamma\delta\beta}}{\partial Q_p} \right) \Big|_{r_e} \quad (18)$$

In these equations, ν_0 and ν_{1p} denote the vibrational ground state and the first vibrationally excited state for vibrational mode p ,

TABLE 1: Calculated ROA Intensities for *R*-Methyloxirane^a

6-31G				6-31G**				cc-pVDZ				aug-cc-pVDZ				experiment	
ω^{HF}	ω^{DFT}	Δ_z^{HF}	Δ_z^{DFT}	ω^{HF}	ω^{DFT}	Δ_z^{HF}	Δ_z^{DFT}	ω^{HF}	ω^{DFT}	Δ_z^{HF}	Δ_z^{DFT}	ω^{HF}	ω^{DFT}	Δ_z^{HF}	Δ_z^{DFT}	ω	Δ_z
1676	1551	1.4	-0.7	1687	1548	5.9	2.8	1665	1517	8.7	4.3	1659	1518	6.2	5.5	1498	1.0
1652	1541	-1.4	-1.2	1625	1510	-4.6	-4.1	1595	1468	-6.1	-5.9	1595	1472	-7.2	-6.6	1460	-2.7
1641	1528	0.5	1.4	1611	1497	1.3	1.6	1580	1453	2.2	2.9	1578	1455	2.2	2.1	1450	0.8
1592	1464	1.8	0.6	1586	1453	4.4	3.7	1563	1427	3.5	3.6	1561	1429	2.6	2.2	1403	3.1
1567	1446	0.9	0.1	1535	1415	-1.2	-1.1	1507	1384	-2.0	-1.9	1504	1383	-1.8	-1.7	1365	<i>b</i>
1389	1288	0.0	0.0	1410	1300	-0.3	-0.5	1397	1284	-0.6	-0.7	1392	1284	-1.0	-0.9	1262	<i>c</i>
1318	1210	5.9	1.1	1316	1191	6.1	6.3	1296	1172	7.3	9.8	1291	1176	8.4	15.1	1164	4.0
1296	1169	-0.6	-0.4	1296	1167	1.1	-5.3	1276	1152	0.1	-3.5	1274	1152	0.5	-1.9	1140	-3.3
1276	1153	-3.3	-3.1	1271	1155	-8.7	2.0	1259	1138	-8.6	1.9	1258	1143	-7.3	-1.0	1135	-3.8
1218	1117	-1.5	-1.6	1240	1135	-1.9	-2.5	1223	1114	-2.4	-1.8	1217	1112	-2.7	0.6	1101	1.2
1139	1035	20.6	15.5	1141	1042	2.1	-1.4	1128	1027	0.1	-4.0	1123	1027	-1.3	-5.4	1020	-5.1
1049	969	-5.6	-4.3	1077	983	-2.3	-1.5	1071	974	-2.1	-0.7	1059	965	-2.2	-1.9	946	8.3
995	909	2.8	3.3	990	909	2.6	2.4	979	898	2.0	2.1	973	896	2.6	3.2	892	4.6
872	805	1.7	1.8	951	856	1.4	2.6	941	848	1.7	3.1	929	837	1.5	2.6	824	1.5
806	713	-0.6	-1.4	858	782	0.2	-1.5	852	774	0.2	-1.8	844	762	0.1	-2.0	742	-3.1
440	408	0.0	1.3	442	409	-5.2	-2.5	440	408	-2.3	-0.5	439	410	0.8	2.0	419	5.0 ^d
399	360	5.1	7.2	398	365	5.0	6.9	396	363	4.1	6.1	394	365	2.9	2.9	360	2.5 ^d
210	204	0.2	-1.2	222	211	1.7	1.8	225	207	2.0	1.1	225	202	2.1	-2.4	200	+ ^b

^a The Hartree–Fock results have been taken from ref 7 and the experimental observations from ref 34. The frequencies are reported in cm^{-1} .

^b Weak feature. ^c Possible polarization artifact contribution. ^d From ref 42.

and ω_p is the harmonic frequency of this vibration. The derivatives of the generalized polarizability tensors with respect to displacements along the normal coordinate Q_p is evaluated at the equilibrium geometry r_e .

An analytical approach to the calculation of first-order geometrical derivatives of frequency-dependent polarizabilities was recently presented at the Hartree–Fock level.²⁷ However, because of the use of London-atomic orbitals in our calculations, this formalism cannot easily be extended to the electric dipole–magnetic dipole polarizability. We therefore calculate the geometrical derivatives of the polarizability tensors by numerical differentiation, as previously described.⁷

3. Computational Details

Two sets of calculations are presented in this paper. The first is a study of methyloxirane, where special attention is paid to basis-sets effects, as well as to the importance of electron correlation by comparison with previous Hartree–Fock results.⁷ Two types of basis sets are used: the 6-31G and 6-31G** basis sets by Pople and co-workers^{28–31} and the (augmented) correlation-consistent polarized valence double- ζ (aug)-cc-pVDZ basis sets of Dunning and Woon.^{32,33} For comparison, experimental data are included from ref 34.

The second study focuses on α -pinene and *trans*-pinane. For these molecules, we compare both with recent experimental observations¹⁸ and with the SOS approach of Bouř at the DFT level.¹⁷ Due to computational limitations, we have restricted ourselves to using the 6-31G and 6-31G** basis sets for these molecules.

In all calculations, we have used a step length of $0.001a_0$ in the numerical differentiation, somewhat longer than we have used previously.⁷ However, test calculations indicated no change in the calculated CIDs. The numerical differentiation was performed by Cartesian displacements of all the atoms, without the use of translational and rotational invariance. The differentiated tensors were transformed to the normal coordinate basis using an analytically calculated Hessian. In all VROA calculations, the density was converged to a gradient of 10^{-7} and the iterative solution of the linear response equations to 5×10^{-5} . We used a dense grid consisting of 64 radial and 590 angular grid points for each atom. The B3LYP functional³⁵ was used

in all calculations, and all CIDs have been calculated for a wavelength of the incident light of 488.8 nm.

For the calculation of the CIDs, we have used a local version of the Dalton program,²⁴ using the recent implementation of linear response functions for the calculation of optical rotation at the DFT level.¹⁹ However, Dalton is currently incapable of calculating analytical Hessians at the DFT level, for which we instead used Gaussian98.³⁶ It is here important to note that the B3LYP functional as defined in Dalton differs from that used in Gaussian98. Whereas the latter uses the functional form III of the Vosko, Wilk, and Nusair (VWN) correlation functional,³⁷ Dalton uses functional form V of the same paper. However, test calculations of optical rotations show only minor differences compared to the work of Stephens et al.,³⁸ and we will assume that this agreement also holds for the vibrational frequencies. The Hessians were transformed to Dalton format using an interface routine written by Hangartner,³⁹ which is distributed with the Dalton program.

4. Results

4.1. Methyloxirane. Our results for *R*-methyloxirane are collected in Table 1, where we also report the Hartree–Fock results of ref 7 and the experimental results of ref 34. Some of the Hartree–Fock results in Table 1 have been obtained by us to investigate more carefully the importance of diffuse functions on the VROA calculation.

The calculated harmonic frequencies listed in Table 1 confirm the established knowledge that DFT and, in particular, the B3LYP functional, gives better agreement with experiment than Hartree–Fock does for the vibrational frequencies.¹³ There are some changes in the frequencies from 6-31G to 6-31G**, but overall the cc-pVDZ basis set provides a significantly better agreement with experiment than do the Pople basis sets. However, as we compare experimental fundamental frequencies with theoretical harmonic frequencies (and ignore possible solvent effects), exact agreement cannot be expected. The differences between the cc-pVDZ and aug-cc-pVDZ frequencies are negligible. This suggests that augmentation has little influence on the vibrational frequencies.

Turning our attention to the calculated VROA CIDs for right-angle scattering, we note that Hartree–Fock and DFT are in

much better agreement than was the case for the frequencies. Although there are still substantial differences, most of the qualitative intensity features of the spectrum are reproduced by both Hartree–Fock theory and DFT. This observation suggests that the most important aspect for the prediction of a VROA spectrum is an accurate force field and that it should be possible to use the DFT/B3LYP force field together with Hartree–Fock intensities, as done for instance in ref 15. We shall return to this point shortly.

Despite the similarities between the HF and DFT intensities, there are some important differences. In particular, the HF and DFT intensities differ in sign for several of the observed transitions. For the larger basis sets, these differences are mostly confined to low-intensity vibrational modes, thus possibly arising from a partial cancellation of the three contributing terms to the VROA intensity. Still, the differences between HF and DFT are to some extent less severe than the errors incurred by using small basis sets.

Whereas small differences are observed between 6-31G** and cc-pVDZ for the frequencies, larger differences are found for the intensities. It is discouraging to note that there are also significant changes in the intensities from cc-pVDZ to aug-cc-pVDZ. In particular, at the DFT/aug-cc-pVDZ level the number of modes calculated with incorrect signs decreases to two (those at 965 and 202 cm^{-1}), whereas for most other models at least four differences in sign can be observed. Fortunately, only three bands differ in sign at the HF/6-31G level.

In the first theoretical investigation of *R*-methyloxirane, Bose et al. suggested interchanging the bands at 1218 and 1139 cm^{-1} , as these bands at the HF/6-31G level were predicted to be very strong with the opposite sign of experiment. However, the intensity of these bands were reduced significantly in the 6-31G* basis,³⁴ indicating that basis-set effects may be substantial. Indeed, this has been corroborated by the use of the even larger aug-cc-pVDZ basis in ref 7, where both bands become negative. The results presented in Table 1 for these bands suggest that proper treatment of electron correlation (best represented at the B3LYP/aug-cc-pVDZ level) is necessary for a reliable computation of the signs and relative intensities.

Although the cost of the B3LYP calculation scales similarly as for the HF calculation, the DFT method involves a numerical quadrature in addition to the Coulomb and exchange integrals, thus making the calculation always more expensive. The cost is further increased by the need of a fine integration grid to avoid problems with inaccuracies in the numerical differentiation. It is therefore of practical interest to explore the possibility of combining the DFT force field with the HF polarizability derivatives, as in Table 2, where we report the results obtained by evaluating the HF ROA intensities both at the HF and B3LYP minimum geometry, with the aug-cc-pVDZ basis. Although pure HF and B3LYP computations differ significantly (see Table 1), the combination of the B3LYP force field and the HF CIDs provides results of the same quality as those obtained with pure DFT. Moreover, there are no significant differences between CIDs evaluated at the HF and B3LYP minima. Because, however, the force field needs to be determined at the B3LYP minimum, the DFT geometry appears more convenient.

There are two minor exceptions to the agreement between DFT and the mixed DFT/HF scheme. For the frequency at 1135 cm^{-1} , use of the Hartree–Fock geometry in the pure DFT scheme produces an intensity with a sign different from that of the other methods and that of experiment. The situation is different for the mode at 200 cm^{-1} , where only the intensities

TABLE 2: Calculated ROA Intensities for *R*-Methyloxirane^a

ω_{exp}	HF geometry		DFT geometry		Δ_z^{exp}
	Δ_z^{HF}	Δ_z^{DFT}	Δ_z^{HF}	Δ_z^{DFT}	
1498	6.2	6.1	6.2	5.5	1.0
1460	-7.2	-6.3	-6.6	-6.6	-2.7
1450	2.2	2.1	2.0	2.1	0.8
1403	2.6	2.5	2.5	2.2	3.1
1365	-1.8	-2.5	-2.1	-1.7	<i>b</i>
1262	-1.0	-1.0	-1.0	-0.9	<i>c</i>
1164	8.4	17.0	16.6	15.1	4.0
1140	0.5	-1.8	-1.8	-1.9	-3.3
1135	-7.3	0.2	-0.3	-1.0	-3.8
1101	-2.7	-2.2	-1.6	0.6	1.2
1020	-1.3	-3.6	-4.2	-5.4	-5.1
946	-2.2	-2.8	-2.9	-1.9	8.3
892	2.6	3.2	3.0	3.2	4.6
824	1.5	1.9	1.9	2.6	1.5
742	0.1	-0.6	-0.6	-2.0	-3.1
419	0.8	1.0	1.1	2.0	5.0 ^d
360	2.9	3.1	3.2	2.9	2.5 ^d
200	2.1	-1.9	1.2	-2.4	+ ^b

^a The Hartree–Fock results have been taken from ref 7 and the experimental observations from ref 34. The frequencies are reported in cm^{-1} . In all cases a DFT/B3LYP force field has been used. The superscript in Δ_z^{wf} indicates the wave function used for calculating the VROA intensities. ^b Weak feature. ^c Possible polarization artifact contribution. ^d From ref 42.

calculated in the mixed approach has the same sign as that observed in experiment. Because of the flat potential for this mode, which involves large relative atomic displacements as it is dominated by methyl rotation, small differences in the equilibrium geometry may give large differences in the VROA intensities. We therefore cannot draw definite conclusions for this vibrational mode. Also, because the internal rotation is almost free, the perturbational treatment used here for calculating the VROA intensities cannot be expected to be accurate.

4.2. α -Pinene and *trans*-Pinane. Our calculated B3LYP frequencies and VROA circular intensity differences $I_{180}^{\text{R}} - I_{180}^{\text{L}}$ for (1*S*)-(–)- α -pinene in the 6-31G and 6-31G** basis sets are collected in Table 3, along with the results of the sum-over-states approach and the mixed HF/DFT scheme of Bouř.¹⁷ The force field and geometry of all calculations in ref 17 were obtained using the BPW91 functional and a TZ2P basis; their results are therefore not directly comparable with ours, obtained with the B3LYP functional and the 6-31G and 6-31G** basis sets. However, the SOS CIDs of ref 17 were obtained at the B3LYP/6-31G** level. The mixed HF/DFT calculations employed Hartree–Fock intensities using the 6-31G** basis.

There are significant differences in the frequencies at the BPW91/TZ2P and B3LYP/6-31G** levels, whereas the B3LYP frequencies are rather similar for the 6-31G and the 6-31G** basis sets. In view of the discussion in the previous section, the main reason for the differences is probably the use of a much larger basis for the calculation of the molecular force field in refs 17 and 40.

Considering the differences in the force field, the SOS scheme performs surprisingly well compared to our full DFT calculation. The main features of the spectrum are in good agreement with the B3LYP/6-31G** intensities, although the magnitude of the CIDs are often overestimated. For a few bands, the different models predict different signs of the CIDs, in particular in the 1100–900 cm^{-1} region. For most bands in this region, the HF/DFT results of ref 17 and our B3LYP/6-31G results agree with the “reference” B3LYP/6-31G** results, indicating that the SOS+Polar scheme does not reproduce the CIDs correctly in this region.

TABLE 3: Vibrational Frequencies (in cm^{-1}) and Backscattering CIDs of α -Pinene

ref 17	vibrational frequency		CIDs			
	B3LYP/6-31G	B3LYP/6-31G**	Polar+SOS ¹⁷	mixed HF/DFT ¹⁷	B3LYP/6-31G	B3LYP/6-31G**
1663	1734	1732	-34	-12	-8	-9
1477	1561	1531	2	11	16	8
1468	1555	1526	-62	-71	-34	-25
1455	1541	1511	-19	-32	-48	-44
1451	1538	1505	-195	-257	-56	-103
1448	1534	1502	307	413	281	256
1440	1527	1496	77	100	32	45
1437	1524	1490	4	8	24	21
1432	1524	1490	-174	-255	-276	-217
1373	1463	1431	3	9	8	7
1365	1460	1426	9	18	7	13
1354	1445	1413	5	1	6	4
1326	1394	1375	-3	0	5	5
1312	1382	1366	-4	6	9	15
1292	1358	1345	-5	-8	5	-6
1250	1311	1305	13	10	-12	1
1235	1304	1283	8	11	4	-1
1206	1274	1258	34	37	25	17
1199	1257	1244	-9	-32	-4	-13
1190	1250	1237	-44	-39	-30	-19
1169	1238	1215	2	0	5	-1
1155	1211	1195	19	25	14	19
1114	1170	1154	-17	-8	-16	-8
1087	1146	1134	21	32	-1	-5
1074	1127	1113	-7	43	49	37
1051	1112	1085	15	2	7	-1
1031	1097	1068	-40	-30	-26	-14
1026	1074	1061	18	-15	-12	-2
1003	1059	1039	-12	-5	-10	0
989	1047	1024	-1	8	-5	-5
944	997	975	-30	-32	12	-12
939	988	975	-29	12	4	16
923	971	954	6	-2	-9	-3
916	953	944	47	60	28	29
893	928	923	33	35	26	18
873	905	901	-8	-16	-23	-19
830	854	856	-15	-25	-11	-8
800	832	829	5	10	-13	1
777	824	813	-65	-95	-34	-44
758	786	783	24	45	21	20
655	680	674	5	17	6	7
605	636	626	83	60	31	25
556	581	574	-55	-48	-22	-17
473	495	486	3	11	7	7
451	473	465	18	-5	-1	-1
410	433	429	-8	-3	-3	-4

In Figure 1, we have plotted the experimental spectrum of (1*S*)-(–)- α -pinene from ref 18 together with a simulated spectrum based on our B3LYP/6-31G** results, using a Lorentzian line width of 8 cm^{-1} . We note that the experimental spectrum has a rather low signal-to-noise (S/N) ratio, and that VROA spectra of this molecule, as well as *trans*-pinane, with a much improved S/N ratio have been published by Vargek, Freedman, and Nafie.⁴¹ However, despite the low S/N ratio, it can easily be seen that the simulated spectrum is in good agreement with the experimental spectrum we have plotted in Figure 1.

In Figure 2, we have plotted the experimental Raman spectrum of (1*S*)-(–)- α -pinene together with our simulated spectrum obtained using B3LYP/6-31G**. For the simulated Raman spectra, we have plotted the absolute differential Raman scattering cross section,⁹ setting the height of the peak at 1450 cm^{-1} to approximately the same height as in the experimental Raman spectrum. Although many of the qualitative features of the Raman spectrum are correctly reproduced, we note that there

TABLE 4: Vibrational Frequencies (in cm^{-1}) and Backscattering CIDs of *trans*-Pinane

ref 17	vibrational frequency		CIDs		
	B3LYP/6-31G	B3LYP/6-31G**	Polar+SOS ¹⁷	B3LYP/6-31G	B3LYP/6-31G**
1568	1537		-41	-58	-50
1561	1530		22	51	24
1549	1518		84	68	55
1548	1517		-39	-23	-3
1542	1512		47	-50	42
1542	1510		-53	39	-47
1538	1504		-22	-5	11
1532	1498		-67	3	-76
1529	1497		46	-44	25
1464	1433		6	4	3
1454	1422		2	-4	-1
1446	1415		1	2	4
1412	1397		10	47	18
1401	1379		-7	-41	3
1383	1362		-36	-5	-29
1377	1351		14	22	21
1346	1337		66	55	50
1340	1324		-107	-83	-92
1316	1304		-7	-11	13
1290	1273		14	22	14
1271	1254		-3	-14	-12
1258	1242		8	15	4
1253	1235		-13	-4	-12
1241	1229		-6	-15	1
1211	1195		-25	-5	-49
1197	1182		-1	5	7
1147	1127		57	60	57
1128	1112		26	33	19
1122	1102		-25	-41	-22
1088	1076		-18	-27	-16
1070	1057		6	-19	-15
1048	1027		28	19	24
1029	1014		-27	14	6
1026	1012		25	23	8
988	969		-76	-21	-62
981	965		19	-4	31
968	952		-4	-8	-4
937	933		95	56	61
916	907		-18	-6	-11
895	890		67	34	43
863	864		-141	-105	-93
832	830		2	-3	0
799	789		32	47	39
790	785		-18	-1	-6
668	664		0	1	-5
586	576		1	-4	-4
533	522		5	9	7
487	480		-7	-6	-6
460	452		1	1	1
436	431		0	0	0

are significant discrepancies between the experimental Raman spectrum and our simulated spectrum, in particular in the regions $1400\text{--}1200$ and $950\text{--}850 \text{ cm}^{-1}$. This might indicate a need for more diffuse functions to properly describe the polarizability derivatives but may also be due to solvent effects (not included in our study) or conformational flexibility of the molecule. Considering these problems for the Raman spectrum, the agreement observed for the VROA spectrum using the 6-31G** basis set is promising.

In Table 4, we have collected our results for the frequencies and CIDs in (–)-*trans*-pinane. For this molecule, we have also calculated the CIDs using the SOS+Polar model of ref 17 at the B3LYP/6-31G** level, enabling a direct comparison of the SOS+Polar results with the full DFT results.

In view of its low computational cost, the SOS+Polar model performs very well. Interestingly, for many of the CIDs where

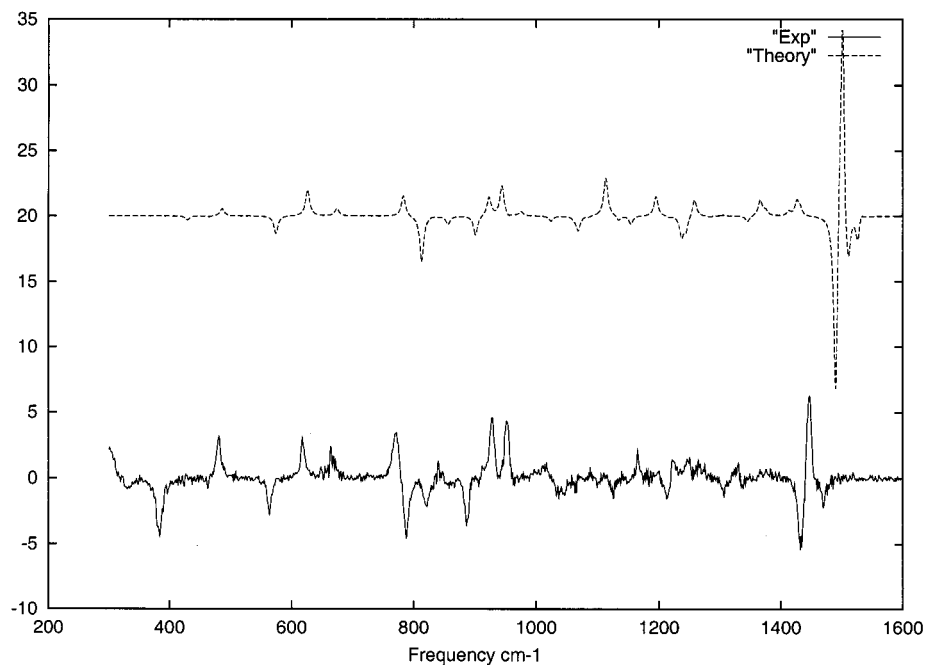


Figure 1. Experimental (from ref 18) and simulated spectra of VROA backscattering of α -pinene. Units on the y axis are arbitrary. A line width of 8 cm^{-1} was used in the simulation of the theoretical spectrum.

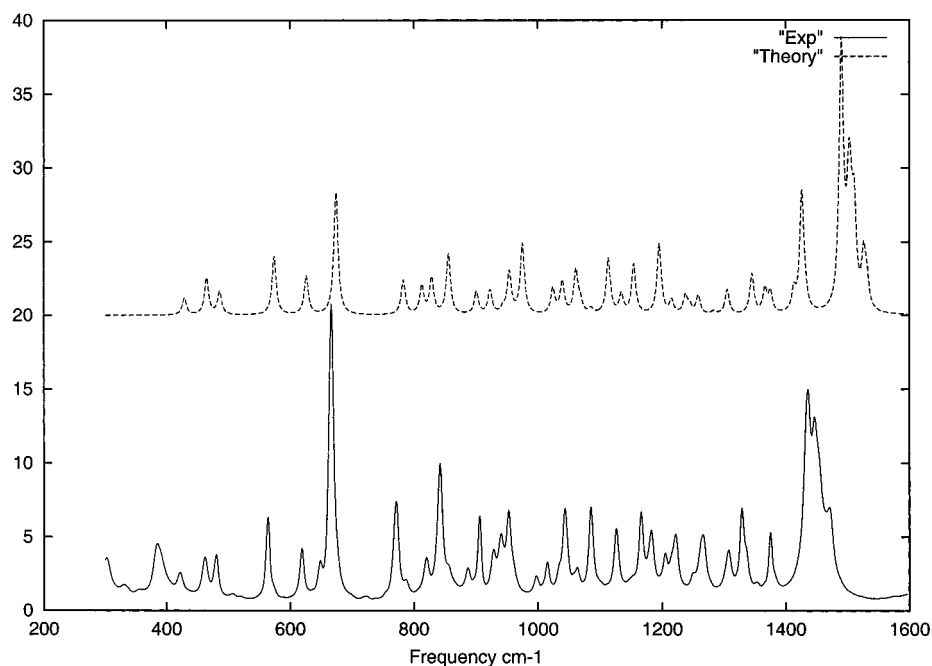


Figure 2. Experimental (from ref 18) and simulated Raman spectra of α -pinene. Units on the y axis are arbitrary. A line width of 8 cm^{-1} was used in the simulation of the theoretical spectrum.

the SOS+Polar model differs from the DFT/6-31G** approach, it agrees more with B3LYP/6-31G. In some sense, therefore, the SOS model mimicks the effect of using smaller basis sets, presumably leading to a poorer description of the excited states. In general, the SOS+Polar model works better for vibrational modes with large intensity differences, where it often gives better agreement with the full B3LYP/6-31G** results than with the B3LYP/6-31G results.

In Figure 3, we have plotted a simulation of our B3LYP/6-31G** and SOS+Polar results together with the experimental VROA spectrum of ref 18. As for α -pinene, the experimental spectrum is well reproduced by the B3LYP/6-31G** calculation, and also by the SOS+Polar approach

thus appears to be a fast and inexpensive tool for quickly providing theoretical CIDs that may help in the first assignment of the experimental VROA spectra. On the other hand, the consistent linear response DFT ROA computation is clearly necessary to reproduce most of the finer spectral features. These include the negative signal around 650 cm^{-1} missing for the SOS model, as well as the small positive bands around 940 and 1140 cm^{-1} , or the positive signal around 1330 cm^{-1} , overestimated by SOS.

In Figure 4 we have plotted the experimental and theoretical Raman spectrum of (*-*)-*trans*-pinane. The simulated Raman spectrum for this molecule seems to be in better agreement with the general intensity features of the experimental spectrum than

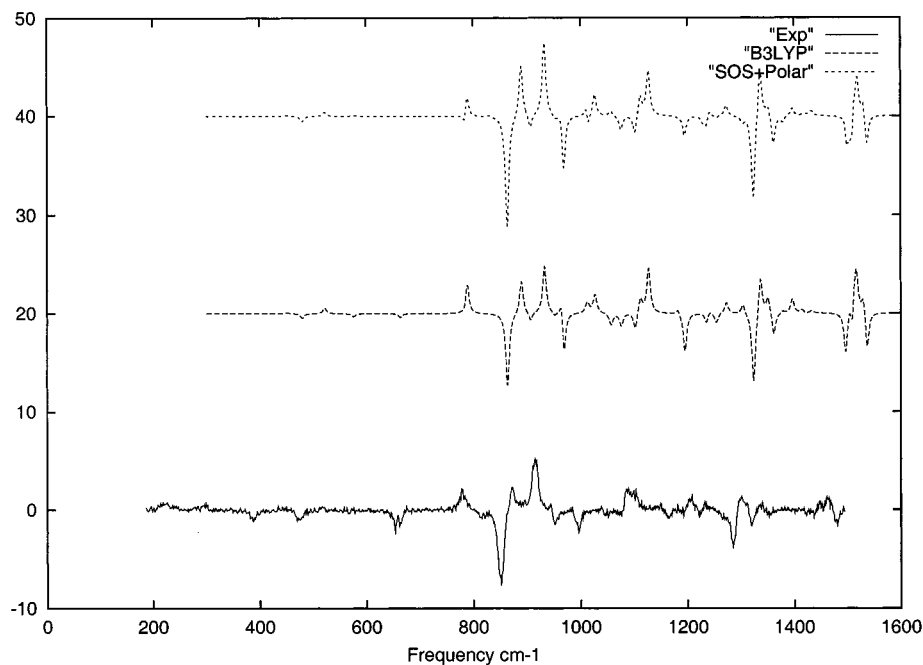


Figure 3. Experimental (from ref 18) and simulated spectra of VROA backscattering of *trans*-pinane. Units on the y axis are arbitrary. A line width of 8 cm^{-1} was used in the simulation of the theoretical spectrum.

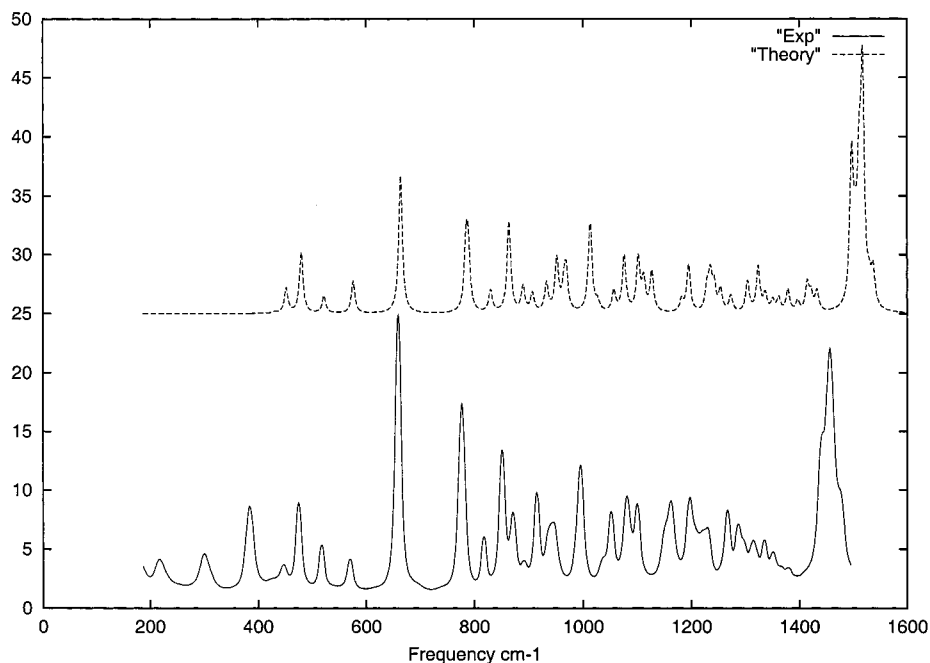


Figure 4. Experimental (from ref 18) and simulated Raman spectra of *trans*-pinane. Units on the y axis are arbitrary. A line width of 8 cm^{-1} was used in the simulation of the theoretical spectrum.

was the case for (1*S*)-(–)- α -pinene. However, in the region $1400\text{--}1200\text{ cm}^{-1}$, the agreement with experiment is rather poor.

5. Concluding Remarks

We have presented the first gauge-origin independent calculations of vibrational Raman optical activity consistently using linear response methods implemented within density-functional theory; the frequency dependence has been included in the calculation, and gauge-origin independence has been ensured by the use of London atomic orbitals. It has been confirmed that use of the DFT/B3LYP force field instead of a Hartree–Fock force field leads to a significant improvement in the

harmonic vibrational frequencies. Although the use of DFT also improves the VROA intensities, the differences between HF and DFT for the CIDs are in general much smaller than for the frequencies. However, to reproduce the finer details of the VROA spectrum, the presented DFT linear response approach appear to be the most accurate method currently available.

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